

# Synthesis of $\alpha$ -MoO<sub>3</sub> nanofibers for enhanced field-emission properties

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## Abstract

One-dimensional  $\alpha$ -MoO<sub>3</sub> nanofibers of 280–320 nm diameters were synthesized by a hydrothermal method. The morphologies and compositions of as-synthesized  $\alpha$ -MoO<sub>3</sub> nanofibers have been characterized by X-ray powder diffraction, Raman spectroscopy, and field-emission scanning electron microscopy. X-ray photoelectron spectroscopy showed the predominantly 6+ oxidation state with a small percentage of reduced  $\delta^+$  ( $5 < \delta < 6$ ) oxidation state. The field-emission properties of  $\alpha$ -MoO<sub>3</sub> nanofibers show a lower turn-on electric field of 2.48 V  $\mu\text{m}^{-1}$  and threshold electric field of 3.10 V  $\mu\text{m}^{-1}$ . The results suggest that the  $\alpha$ -MoO<sub>3</sub> nanofibers are promising candidate for efficient and high-performance field-emission devices. Copyright © 2018 VBRI Press.

**Keywords:** Nanofibers; molybdenum oxide; hydrothermal method; field emission properties.

## Introduction

One-dimensional (1D) nanostructures have been the subject of extensive research over few decades due to their peculiar and fascinating properties with novel applications [1-3]. Moreover, the flexibility, shape anisotropy, extraordinary lengths and high surface to volume ratio of these nanostructures explore the use from electronic to photonic devices [4-8]. Different sizes and morphology of the semiconductor nanostructures are reported for the potential application of field emitters such as field-effect transistors, lasers diodes, sensors, and photovoltaic cells [9-11].

Recently, the  $\alpha$ -MoO<sub>3</sub> one of the utmost stable form of the molybdenum metal oxide semiconductor (band-gap = 3.8 eV), shows a remarkable attention due to extensive potential applications in photoconductive, thermoelectric, pseudocapacitors, Li-ion intercalation-deintercalation and field-emission devices [12,13]. Now a days field-emission, is one of the commercial attention for electronic and display devices. Further 1D nanostructure is one of the prominent features for the field-emission in the viewpoint of application. It has enhanced the interest in the synthesis methods as well as electrochemical properties of 1D semiconductor nanostructures. In the past decades, field-emission properties are investigated from different 1D shapes of MoO<sub>3</sub> [14-17]. Yan *et al.* [14] have demonstrated that MoO<sub>3</sub> nanoflakes fabricated by modified hot plate method showed enhanced electron field-emission and

reduced turn-on fields. Li *et al.* [15] have reported an improved field-emission property from the MoO<sub>3</sub> nanobelts synthesized by infrared irradiation heating a Mo foil in the air. Zhou *et al.* [17] have fabricated MoO<sub>3</sub> nanowires via thermal evaporation method followed by further oxidation, and this nanowire exhibited better field-emission than MoO<sub>3</sub> nanoflakes [14] and nanobelts arrays [15]. It is clear from the above report that the material with sharp tips and high aspect ratio can significantly increase field-emission performance [14-17]. Thus nanofibers are expected to possess a better field-emission property owing to their high aspect ratio, exposed sharp edges and low dimension geometry. It is still challenging task for the enhancement of field-emission properties in  $\alpha$ -MoO<sub>3</sub> to meet viable applications required for flat displays, X-ray sources and development of microwave circuits.

In present study, we report a facile synthesis method for preparing  $\alpha$ -MoO<sub>3</sub> nanofibers. The field-emission properties of the prepared nanofibers were investigated. These  $\alpha$ -MoO<sub>3</sub> nanostructures shows excellent field-emission properties with stable emission behavior and low turn-on electric fields.

## Experimental

The  $\alpha$ -MoO<sub>3</sub> nanofibers were synthesized as per our previously reported method [13]. A freshly prepared 0.7 g polymeric nitrosyl-complex of molybdenum (II) [18] was dissolved in a mixture of 35 mL distilled water and

concentrated  $\text{HNO}_3$  under constant stirring. The resulting yellowish-orange colored solution was transferred into an autoclave and then heated in a furnace at  $220^\circ\text{C}$  for 7 h. The obtained material was cooled to room temperature and thoroughly washed with deionized water and ethanol. The samples were collected and dried at  $60^\circ\text{C}$  for 5 h.

The crystal structures were examined by X-ray diffraction (XRD) using a Rich Seifert Isoflex diffractometer with  $\text{Cu-K}\alpha$  ( $\lambda = 1.5418 \text{ \AA}$ ) radiation. The vibrational modes of  $\alpha\text{-MoO}_3$  were characterized by Raman spectrometer (WITec Alpha\_SNOM,  $\lambda = 532 \text{ nm}$  line laser as excitation source). A field-emission scanning electron microscope (FESEM) FEI Quanta-200 was used to analyze the morphology of the sample. The chemical composition was analyzed using X-ray photoelectron spectroscopy (XPS) (SPECS, Phoibos 100 MCD Energy Analyzer). The field-emission studies were done by ‘close proximity’ (also termed as ‘planar diode’) set up using ultrahigh vacuum ( $\sim 1 \times 10^{-8}$  Torr pressure). As synthesized  $\alpha\text{-MoO}_3$  nanofibers deposited onto tungsten substrate ( $1 \times 1 \text{ cm}^2$ ) by drop-casting method. Tungsten and copper were used as a cathode and an anode, respectively. The separation between cathode and anode was fixed at  $250 \mu\text{m}$  by using alumina spacer. Keithley 485 Pico ammeter was used to carry out for the current density–applied field ( $J$ – $E$ ) measurements.

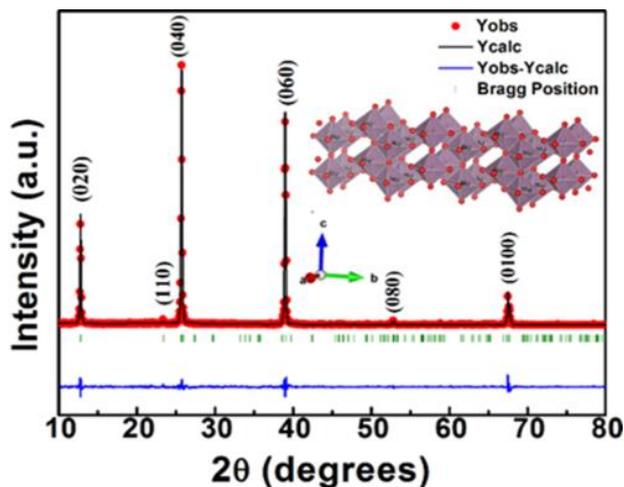


Fig. 1. Rietveld analysis of XRD data using orthorhombic structure with space group Pbnm of  $\alpha\text{-MoO}_3$  nanofibers and the insert shows the unit cell crystal structure.

## Results and discussion

Fig. 1 depicts the Rietveld-refined XRD pattern of  $\alpha\text{-MoO}_3$  powder, which is indexed as the orthorhombic crystal system of  $\text{MoO}_3$  having Pbnm space group (JCPDS card No. 76-1003). In XRD, a collimated X-ray beam with wavelength,  $\lambda$  are incident upon a crystalline sample and is diffracted according to the Bragg's law [19]:

$$2d \sin\theta = n\lambda$$

where  $d$  is the atomic spacing of two planes,  $\theta$  is the diffraction angle and  $n$  is the order of diffraction. The lattice parameters for the orthorhombic crystal system ( $a \neq b \neq c$  and  $\alpha = \beta = \gamma = 90^\circ$ ) are related to the  $d$ -spacing of the  $(hkl)$  planes as given by the equation:

$$\frac{1}{d^2} = \frac{h^2}{a^2} + \frac{k^2}{b^2} + \frac{l^2}{c^2}$$

where  $h$ ,  $k$  and  $l$  are the Miller indices of the planes. The obtained lattice parameters are:  $a = 3.9623$ ,  $b = 13.8576$  and  $c = 3.6962 \text{ \AA}$ , respectively and the unit cell volume is  $202.9508 \text{ \AA}^3$ . Reliability factors for the fit are  $R_p = 13.8\%$ ,  $R_{wp} = 25.6\%$ ,  $R_{exp} = 18.32\%$  and  $\chi^2 = 1.95$ . The XRD pattern shows that all other peaks except for  $(0k0)$  have very weak intensity in contrast to standard data. The observed intensities of reflection peaks of  $(020)$ ,  $(040)$ ,  $(060)$ ,  $(080)$  and  $(0100)$  are very strong, indicating the anisotropic morphology growth and preferential orientation of the  $\alpha\text{-MoO}_3$  nanofibers along the  $[001]$  direction [13].

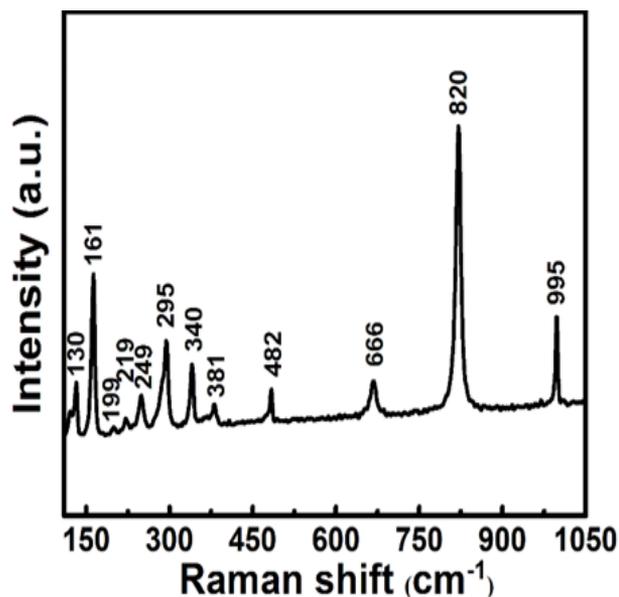


Fig. 2. Raman spectra of the  $\alpha\text{-MoO}_3$  nanofibers.

Micro-Raman spectra of  $\alpha\text{-MoO}_3$  nanofibers were collected under  $532 \text{ nm}$  excitation which is shown in Fig. 2. The observed band positions of obtained  $\alpha\text{-MoO}_3$  nanofibers are shown in Table 1. The Raman modes for  $\alpha\text{-MoO}_3$  can be written as;

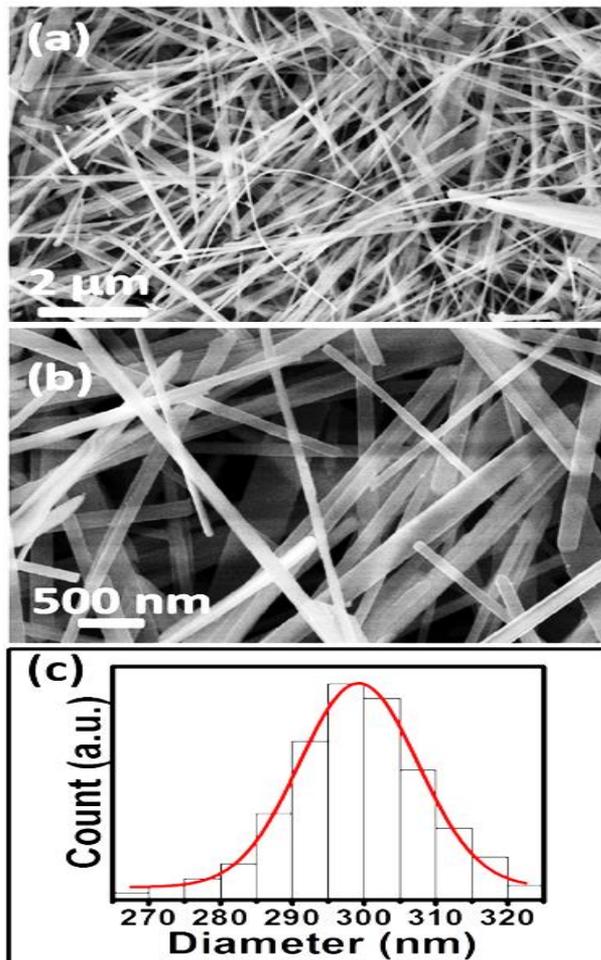
$$\Gamma = 8A_g + 8B_{1g} + 4B_{2g} + 4B_{3g} + 4A_u + 3B_{1u} + 7B_{2u} + 7B_{3u}$$

where  $A_g$ ,  $B_{1g}$ ,  $B_{2g}$ ,  $B_{3g}$  are Raman active,  $A_u$  is an inactive, and others are infrared-active modes [20]. These modes are consistent with those in molybdenum trioxide reported by Lupan et al. [20] as shown in Table 1. Detailed Raman vibrational analysis of  $\alpha\text{-MoO}_3$  bands can be found elsewhere [8,21,22].

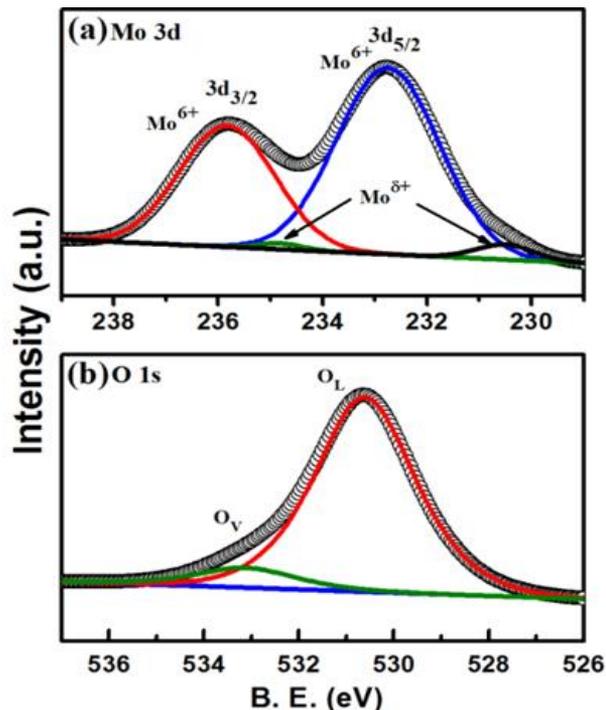
**Table 1.** Peak position for the Raman modes ( $\text{cm}^{-1}$ ) in our study on  $\alpha\text{-MoO}_3$  nanofibers in comparison to the reported data on  $\alpha\text{-MoO}_3$  nano-microribbons by Lupan *et al.* [20]

Raman modes	$\alpha\text{-MoO}_3$ (our study)	Lupan <i>et al.</i>
$B_{2g}$	130	112
$A_g, B_{1g}$	161	154
$B_{2g}$	199	194
$A_g$	219	213
$B_{3g}$	249	-
$B_{2g}, B_{3g}$	295	280
$B_{1g}, A_g$	340	333
$A_g$	381	375
$A_g, B_{1g}$	482	469
$B_{2g}, B_{3g}$	666	662
$A_g, B_{1g}$	820	815
$A_g, B_{1g}$	995	991

**Fig. 3** shows the low and high-resolution FESEM images of  $\alpha\text{-MoO}_3$  nanostructures. As shown in **Fig. 3**,  $\alpha\text{-MoO}_3$  nanostructures have fiber-like shape, which is distinct from those of nanowires and nanotubes, indicating that the hydrothermal method is an efficient method to synthesize the 1D nanostructures. The diameters of the  $\alpha\text{-MoO}_3$  nanofibers are between 280-320 nm and lengths greater than 5  $\mu\text{m}$ .



**Fig. 3.** FESEM images (a) low, (b) high-magnification and (c) histogram for the width distribution of the  $\alpha\text{-MoO}_3$  nanofibers.



**Fig. 4.** XPS spectra of the (a) Mo 3d and (b) O 1s core level in  $\alpha\text{-MoO}_3$  nanofibers.

**Fig. 4** shows the X-ray photoelectron spectroscopy (XPS) core level spectra of  $\alpha\text{-MoO}_3$  nanofibers, calibrated by the C (1s) line (285.0 eV) binding energy. In **Fig. 4(a)** two different valence states of Mo (3d) in the  $\alpha\text{-MoO}_3$  are observed. The major contributing peaks at 232.7 eV ( $3d_{5/2}$ ) and 235.8 eV ( $3d_{3/2}$ ) corresponded to  $\text{Mo}^{6+}$ , while the small ones at 230.5 eV ( $3d_{5/2}$ ) and 234.8 eV ( $3d_{3/2}$ ) are assigned to the 3d doublet of  $\text{Mo}^{\delta+}$  ( $5 < \delta < 6$ ) [8,23]. It indicates the existence of dangling bond defects in  $\alpha\text{-MoO}_3$  where charges could be trapped [8,23]. The mixed valence state of Mo ( $\text{Mo}^{6+}$  and  $\text{Mo}^{\delta+}$ ) is often followed by the introduction of oxygen vacancies to keep charge neutrality [8,23]. Asymmetric and broad peaks observed in the O 1s spectrum, as shown in **Fig. 4(b)**, at 530.6 eV and 533.1 eV correspond to the skeletal oxygen ( $O_L$ ) and oxygen vacancies ( $O_V$ ), respectively [23]. Thus, the XPS result confirms that the  $\alpha\text{-MoO}_3$  possess the oxygen vacancies.

**Fig. 5(a)** depicts typical current density versus applied electric field ( $J$ - $E$ ) plot for the  $\alpha\text{-MoO}_3$  nanofibers. A reproducibly emission current density of  $\sim 15 \text{ nAcm}^{-2}$  is observed with an applied electric field of  $1.76 \text{ V}\mu\text{m}^{-1}$ . It is observed that the current density increases exponentially with increase in electric field and found to be of  $100 \mu\text{Acm}^{-2}$  at an applied field of  $4 \text{ V}\mu\text{m}^{-1}$ . Hence, we could defined the turn on and threshold fields as the required electric fields to generate a current density of 1 and  $10 \mu\text{A cm}^{-2}$ , respectively. Lower the value of turn-on field leads to better field-emission properties of the material. In the present study, the value of turn-on and threshold fields are found to be  $2.48$  and  $3.10 \text{ V}\mu\text{m}^{-1}$ , respectively. It

can be seen from the **Table 2**; these observed values are noticeably lower than the values reported for the MoO<sub>3</sub> nanoflakes [14], nanobelts [15], nanoflowers [16], and aligned nanowires [17].

The dependence of current density over the electric field is expressed by using Fowler–Nordheim (F–N) law [24]:

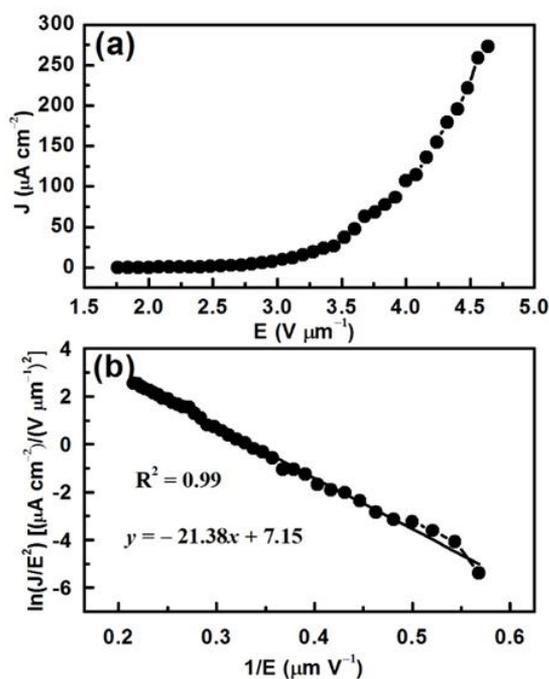
$$J = A \left( \frac{\beta^2 E^2}{\phi} \right) \exp \left( - \frac{B \phi^{3/2}}{\beta E} \right) \text{ or}$$

$$\ln \left( \frac{J}{E^2} \right) = \ln \left( \frac{A \beta^2}{\phi} \right) - \frac{B \phi^{3/2}}{\beta E}$$

The plot of  $\ln(J/E^2)$  versus  $(1/E)$  is termed as F–N plot and derived from experimental data for  $\alpha$ -MoO<sub>3</sub> nanofibers is depicted in **Fig. 5(b)**. The linear nature of F–N plot shows that field-emission from  $\alpha$ -MoO<sub>3</sub> nanofibers follows the quantum mechanical tunneling process [14–17, 25, 26]. The field enhancement factor  $\beta$  is calculated by using the following equation:

$$\beta = \frac{(-6.83 \times 10^3 \phi^{3/2})}{m}$$

where  $\phi$  is the work function of emitted material in eV, which is at least about 5.7 eV for  $\alpha$ -MoO<sub>3</sub> and  $m$  is the slope of the F–N plot [27,28]. Using the slope ( $m = -21.38$ ) of fitted straight line of the F–N plot (**Fig. 5(b)**),  $\beta$  is calculated and found to be 4347, which is very high. The low turn-on field (2.48 V $\mu$ m<sup>-1</sup>) and high field-emission factor ( $\beta = 4347$ ) indicated that the  $\alpha$ -MoO<sub>3</sub> nanofibers could be expected to a wide range of significant applications such as field-emitters and flat-panel display devices [14-17].

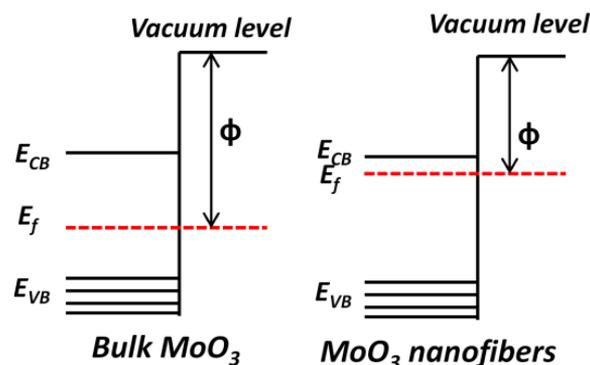


**Fig. 5.** Field-emission properties of  $\alpha$ -MoO<sub>3</sub> nanofibers, (a) current versus electric fields ( $J$ – $E$ ), and (b) corresponding Fowler–Nordheim (F–N) plots with linear curve fitting (solid line).

**Table 2.** Field-emission performances of various  $\alpha$ -MoO<sub>3</sub> nanostructures.

Emitter type	Turn-on field (V $\mu$ m <sup>-1</sup> )	Threshold field (V $\mu$ m <sup>-1</sup> )	$\beta$	Ref.
Nanoflakes	11.0	-	2500	14
Nanobelts	8.7	12.9	-	15
Nanoflowers	4.3	-	-	16
Nanowires	3.5	7.65	4400	17
Nanofibers	2.48	3.10	4347	[this work]

The reasons for our nanofibers possessed excellent field-emission properties may include the following. The low turn-on field and high value of  $\beta$  are attributed to the nanometric features of  $\alpha$ -MoO<sub>3</sub> fibers apex and its high aspect ratio. Also, the defect oxygen vacancies may also play an essential role in describing the surface properties of  $\alpha$ -MoO<sub>3</sub> nanofibers [23]. The presence of oxygen vacancies in  $\alpha$ -MoO<sub>3</sub> nanofibers act as electron donors due to its  $n$ -type semiconducting nature. So these electron donors could shifts the Fermi energy level ( $E_f$ ) up, which leads to lowered work function, as shown in **Fig. 6**. Thus the electron emission process increases due to decrease in the work function of the nanofibers [29]. Also the oxygen ions deficiency give rise to the reduction of the molybdenum ions, that increases the concentration of Mo <sup>$\delta$ +</sup> ( $5 < \delta < 6$ ) ions, which decreases the work function [29,30]. Some further theoretical simulation can be expected to explain this field-emission behavior of  $\alpha$ -MoO<sub>3</sub> nanofibers.



**Fig. 6.** Schematic diagram of energy levels for bulk MoO<sub>3</sub> and MoO<sub>3</sub> nanofibers.  $E_{CB}$ : conduction band.  $E_{VB}$ : valence band.  $E_f$ : Fermi level and  $\phi$ : work function.

## Conclusion

$\alpha$ -MoO<sub>3</sub> nanofibers having an orthorhombic crystal structure were synthesized by a hydrothermal method. The nanofibers are found to be of 280-300 nm in diameter and five to several  $\mu$ m in length. An extremely lower turn-on field (2.48 V  $\mu$ m<sup>-1</sup>), smaller threshold field (3.10 V  $\mu$ m<sup>-1</sup>) and high field-emission factor

( $\beta = 4347$ ) were obtained. Defect mediated (oxygen vacancies) field-emission was proposed to explain this phenomenon. This high-performance emission of  $\alpha$ -MoO<sub>3</sub> nanofibers compared to other MoO<sub>3</sub> nanostructures is controllable, reproducible and can be applied for the development of practical electron sources in flat-panel display devices.

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